

Parametrized tests of post-Newtonian theory with Einstein Telescope

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Tests of General Relativity

Most important Tests till date: [Will, 2001 for a review]

- **Weak-field regime** using **Solar system observations**.
 - ▶ Use of **parametrized post-Newtonian (PPN) formalism**.
 - ▶ Deviation of a general metric theory of gravity in the weak field limit from Newtonian theory was parametrized in terms of 10 free parameters
 - ▶ These parameters are constrained to very good accuracy with various solar system observations.
- **Strong field & Radiative regime** using **binary pulsar observations**:
 - ▶ Strong fields involving compact objects of $v \sim 10^{-3}c$.
 - ▶ Use of **parametrized post-Keplerian (PPK) formalism** as applied to timing equation.
 - ▶ Various Keplerian & post-Keplerian parameters are functions of the individual masses of the binary and determination of more than 2 of these \Rightarrow consistency tests in the $m_1 - m_2$ plane.
- General relativity passes these tests in flying colours!

These tests were so successful because of solid theoretical platforms of PPN and PPK formalisms.

Stretching general relativity further: Tests with GWs

What if

- * General relativity breaks down when the gravitational fields are stronger than those of binary pulsars.
- * There is a scalar field coupled with the metric? [**Scalar-tensor field theories**]
- * Graviton has a mass which is so small that it starts to show up in the very strong field regime. [**Massive Graviton Theories**]
- Gravity is described by some other theory.

Gravitational Waves

- Gravitational Waves have direct imprints of all the strong field effects
- How well can GW observations constrain deviations from GR?

Inspiralling compact binaries and testing general relativity

Adiabatic inspiral phase of a compact binary coalescence is well modelled using post-Newtonian (PN) formalism.

- Determination of coefficients in phasing formula can lead to meaningful tests
 - ▶ Detectability of tails [Blanchet & Sathyaprakash, 1994].
 - ▶ Measuring the dipolar content of the gravitational wave and test scalar-tensor theories [Will, 1994; Krolak et al, 1995, Damour & Esposito-Farèse, 1998].
 - ▶ Parametrizing the 1PN coefficient of the phasing formula capturing the Compton wavelength of the massive graviton and bounding its value from GW observations [Will, 1998].

The question

Can these tests be generalized, without having to know a priori the parameters of the underlying theory of gravity?

Parametrized test of PN theory

Phasing formula in the restricted waveform approximation

$$\tilde{h}(f) = \frac{1}{\sqrt{30} \pi^{2/3}} \frac{\mathcal{M}^{5/6}}{D_L} f^{-7/6} e^{i\psi(f)},$$

and to 3.5PN order the phase of the Fourier domain waveform is given by

$$\psi(f) = 2\pi f t_c - \phi_c - \frac{\pi}{4} + \sum_{k=0}^7 (\psi_k + \psi_{kl} \ln f) f^{\frac{k-5}{3}},$$

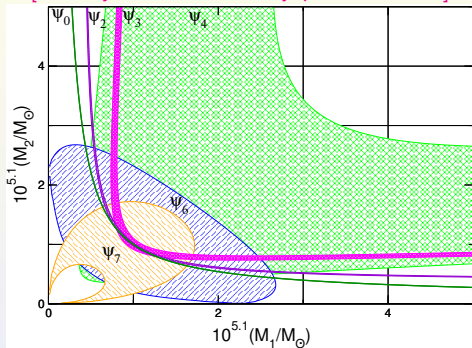
Log terms in the PN expansion

- Phasing coefficients are functions of component masses of the binary: $\psi_k(m_1, m_2)$ & $\psi_{kl}(m_1, m_2)$ [Spins negligible]
- Independent determination of 3 or more of the phasing coefficients \Rightarrow Tests of PN theory [KGA, Iyer, Qusailah & Sathyaprakash, 2006].

Basic Idea

- Parametrize the phasing formula in terms of various phasing coefficients where all of them are treated as independent.
- See how well can different parameters be extracted.
- Those which are well estimated, plot them (ψ_k & ψ_{kl}) in the $m_1 - m_2$ plane (similar to binary pulsar tests) with the widths of various curves proportional to $1 - \sigma$ error bars.

[KGA, Iyer, Qusailah, Sathyaprakash, 2006a]



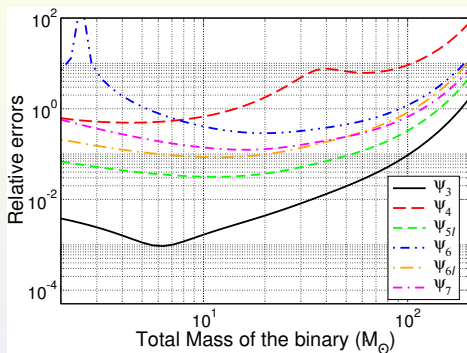
Issues

Highly correlated parameters & Ill-conditioned Fisher matrix for a large parameter space.

Alternative Proposal

[KGA, Iyer, Qusailah & Sathyaprakash, 2006b]

- Treat two parameters as basic variables in terms of which one can parametrize all other parameters EXCEPT one which is the *test* parameter.
- This way, dimensionality of the parameter space is considerably reduced.
- Thus, one will have 8C_3 tests, not all of them independent.
- The best choice to be used as basic variables are the leading two coefficients at 0PN & 1PN, which are the best determined ones.
- Then one will have 6 tests.



- Used an earlier EGO noise PSD (similar to one of the ET noise PSDs).
- All parameters except ψ_4 determined quite well over a large range of masses.

Present work

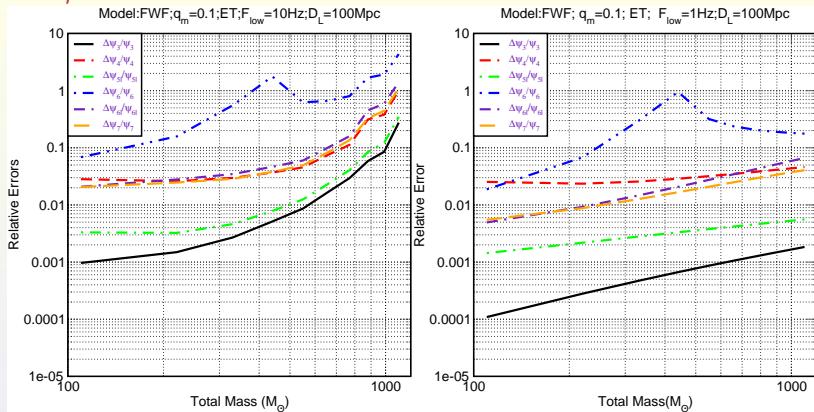
Use of Full Waveforms

- Revisit the earlier estimates more carefully using the ET noise PSD.
- Use of 3PN accurate amplitude corrected waveforms (as opposed to restricted waveforms).
- Effect of low frequency sensitivity on the Test of GR.
- Consideration of unequal mass systems.

Details

- We parametrize the mass dependences (through $\delta = \frac{|m_1 - m_2|}{(m_1 + m_2)}$ and ν) in the *amplitude* terms by ψ_0 & ψ_2 which are used as the basic variables to parametrize all phasing coefficients except the one to be used as test parameter.
- Final parameter space is spanned by: $\{\psi_1, \psi_2, \psi_T, t_c, \phi_c\}$

Results, FWF: 10Hz Cut-off Vs 1Hz Cut-off



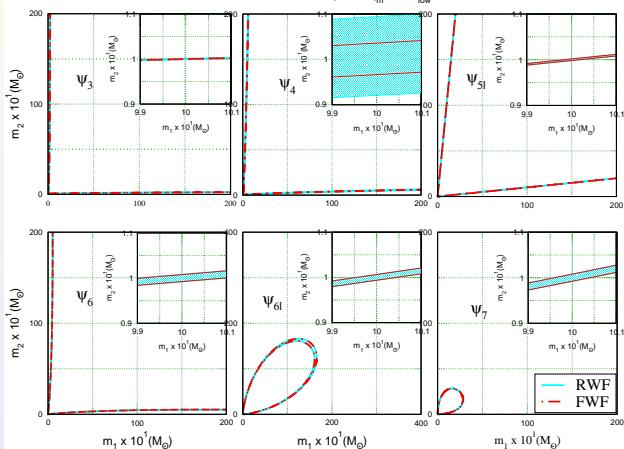
Features

- ψ_3 makes the best use of lowered seismic cut-off.
- Improvement in other parameters is less dramatic due to lower seismic cut-off.
- Improvements are significant for masses $> 500M_\odot$.

$m_1 - m_2$ plots: 1Hz Cut-off (RWF Vs FWF)

System: 10 – 100 M_\odot at 100 Mpc.

RWF vs FWF in m_1 - m_2 plane; $q_m=0.1$; $F_{\text{low}}=1\text{Hz}$



In short

- For this system, FWF doesn't affect the $m_1 - m_2$ plots except for ψ_4 .
- Best test parameter is ψ_3
- Worst test parameter is ψ_4 .

Still..

- Always use the FWF to avoid systematic errors.

Limitations of the proposed Test

- This proposal can test the **overall** consistency between various PN coefficients, **but cannot pin point the inconsistent parameter** as opposed to the earlier proposal which determines all the PN coefficients independently (but is not feasible due to large correlations).
- When different parameters are used as tests, though they are *independent* tests in principle, interpretation of the outcome in the $m_1 - m_2$ plane cannot yield more information other than the overall consistency.
- Cannot be used for probing very specific aspects such as logarithmic terms in the phasing etc.

Other closely related works

- Extracting the three- and four-graviton vertices from binary pulsars and coalescing binaries., [Cannella et al, 2009](#):
 - ▶ Interpreting our tests as measurement of three and four graviton vertices.
- Fundamental Theoretical Bias in Gravitational Wave Astrophysics and the Parametrized Post-Einsteinian Framework., [Yunes & Pretorius, 2009](#):
 - ▶ Biases in using GR templates for GW detection problem and implications.

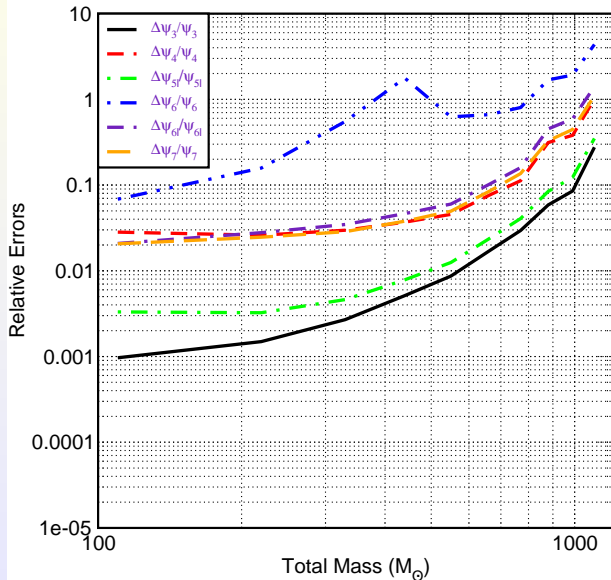
Conclusions & Future directions

- A general way to test the consistency between various PN coefficients, **without having to know the details of the underlying theory**, is proposed.
- Einstein Telescope may have good sensitivity to test the consistency between various PN coefficients in the GW phasing.
- For masses for which the proposed test can be carried out effectively, use of FWF does not bring in dramatic improvement unless ψ_4 is used as test parameter (which is the worst determined parameter when used as test).
(We plan to study this in more detail in the entire parameter space).
- **Still it is strongly advised to use the FWF in these tests to avoid systematic biases due to incomplete waveforms.**
- Lowering the seismic cut-off from 10Hz to 1Hz brings an order of magnitude improvement in the estimation of ψ_3 , which is the best test parameter.
(Currently lookin into how its effect in the $m_1 - m_2$ plane.)
- Effects of **spin**, **residual orbital eccentricity** and the **merger + ringdown** part of the waveforms, on this test have to be investigated (In progress).

Back up slides

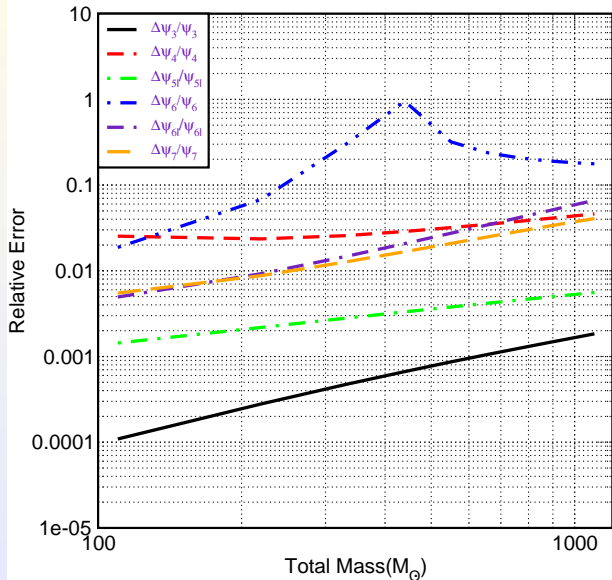
FWF: 10Hz

Model: FWF; $q_m=0.1$; ET; $F_{\text{low}}=10\text{Hz}$; $D_L=100\text{Mpc}$



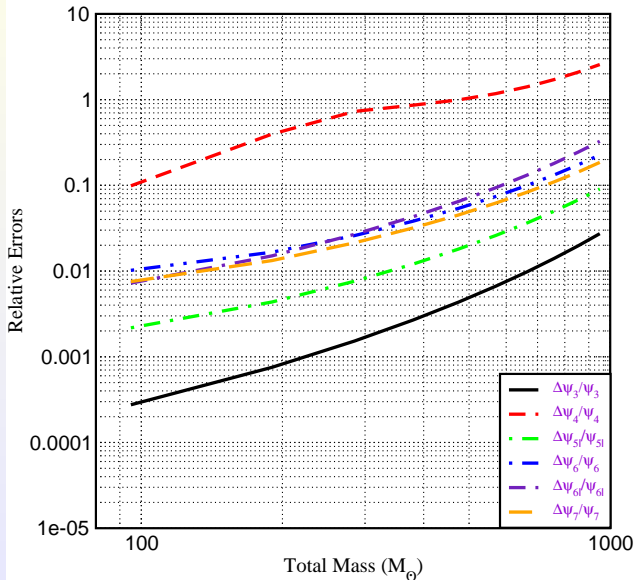
FWF: 1Hz

Model:FWF; $q_m=0.1$; ET; $F_{\text{low}}=1\text{Hz}$; $D_L=100\text{Mpc}$



RWF: 1Hz

Model: RWF; $q_m=0.9$; ET; $F_{\text{low}}=1\text{Hz}$; $D_L=100\text{Mpc}$



FWF: 1Hz

Model: FWF; $q_m=0.9$; ET; $F_{\text{low}}=1\text{Hz}$; $D_L=100\text{Mpc}$

